

# Momentum space analysis of the Triple Pomeron Vertex

Krzysztof Kutak  
Hamburg University  
Institute of Nuclear Physics Kraków

in collaboration with  
Jochen Bartels

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# 1. Motivations



- One of **basic** elements of effective field theory based on **BFKL** pomeron
- Understanding of fan diagram equation for **gluon density**
- Some aspects of the of the pomeron loop

## 2.Two gluons → four gluons effective transition vertex.

(Batels,Wusthoff) Let us consider:

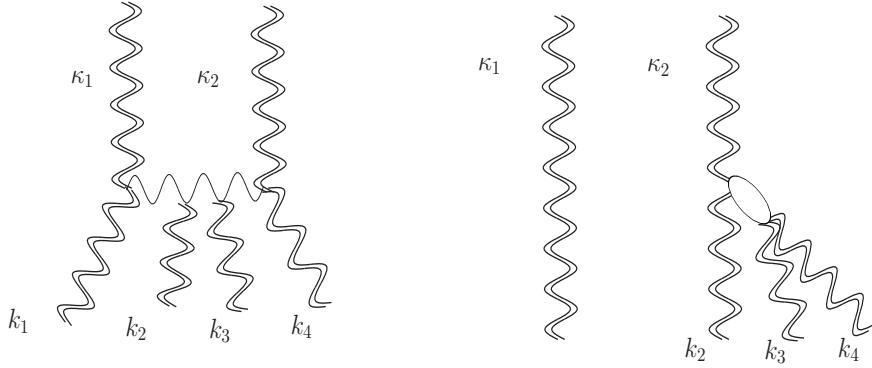


Figure 1: Example of diagrams that correspond to the Vertex

$$\begin{aligned} \mathcal{V}_{a'_1 a'_2; a_1 a_2 a_3 a_4}(\kappa_1, \kappa_2; k_1, k_2, k_3, k_4) = & \\ \frac{\sqrt{2}\pi\delta_{a'_1 a'_2}}{\sqrt{N_c^2 - 1}} & \left[ \delta_{a_1 a_2} \delta_{a_3 a_4} V(1, 2, 3, 4) + \delta_{a_1 a_3} \delta_{a_2 a_4} V(1, 3, 2, 4) \right. \\ & \left. + \delta_{a_1 a_4} \delta_{a_2 a_3} V(1, 4, 2, 3) \right] \end{aligned} \quad (1)$$

where

$$V(1, 2, 3, 4) \equiv V(\kappa_1, \kappa_2; k_1, k_2, k_3, k_4), \quad (2)$$

$$V(1, 2, 3, 4) = \frac{1}{2} g^4 [G(1, 2 + 3, 4) + \dots] \quad (3)$$

$$G(a, b, c) = G_{real}(a, b, c) + G_{virtual}(a, b, c)$$

and  $G_{real} = K_{2 \rightarrow 3}$

In forward case and when  $c = -a$

$$G(a, -, -a) \xrightarrow{\text{red}} K_{BFL}(a, -a)$$

Vertex satisfies Ward identities i.e.  $\mathcal{V}(k_1, k_2, k_3, k_4) \rightarrow 0$  if any  $k_i = 0$

### 3. Vertex in the forward direction

Let us consider configuration:

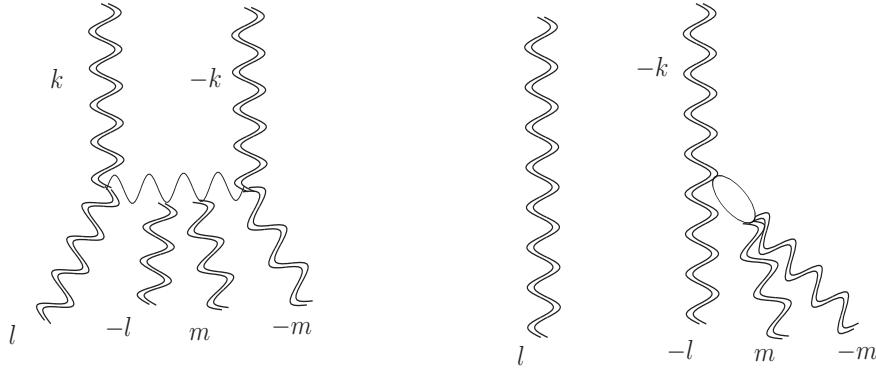


Figure 2:

Vertex simplifies:

$$\begin{aligned}
 \mathcal{V}_{a'_1 a'_2; a_1 a_2 a_3 a_4}(k, -k; l, -l, m, -m) &= \frac{\sqrt{2}\pi\delta_{a'_1 a'_2}}{\sqrt{N_c^2 - 1}} \\
 &\times \left[ \delta_{a_1 a_2} \delta_{a_3 a_4} \left( G(l, -m) + G(-l, -m) + G(l, m) + G(-l, m) \right) \right. \\
 &+ \delta_{a_1 a_3} \delta_{a_2 a_4} \left( G(l, -l) + G(m, -m) + G(-l, m) + G(l, m) - G(l, -l - m) \right. \\
 &\quad \left. \left. - G(-l, l + m) - G(m, -l - m) - G(-m, l + m) + G(l + m, -l - m) \right) \right. \\
 &\quad \left. + \delta_{a_1 a_4} \delta_{a_2 a_3} (G(l, -l) + \dots) \right]
 \end{aligned}$$

goal



twist structure of TPV.

## 4. The Triple Pomeron Vertex Collinear limit

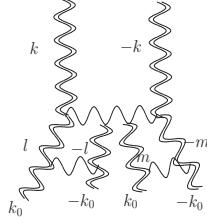


Figure 3:

### Connected part

$$G_1(l, m) = \frac{l^2 k^2}{(k-m)^2} + \frac{m^2 k^2}{(k+m)^2} - \frac{(l+m)^2 k^4}{(k-l)^2 (k+m)^2} \quad (4)$$

$|k| \gg |l|, |m|$  Expanding  $G(l, m)$  in  $l/k, m/k$  we identify:

**twist four**

$$G_1(l, m) \simeq 4\pi k^2 \left[ \frac{2(l \cdot m)^2 - l^2 m^2}{k^4} \right] \quad (5)$$

Projection on color singlets in  $\delta_{a_1 a_2} \delta_{a_3 a_4}$

$$\mathcal{V}_{rl\{a'\}}^{LO N_c}(l, -l, m, -m) = \tilde{\delta}_{a'_1 a'_2} 16\pi \frac{g^4}{2} k^2 \frac{(2(l \cdot m)^2 - l^2 m^2)}{k^4} \quad (6)$$

$$K_{BFKL}^l = \frac{2k_0^2 l^2}{|k_0^2 - l^2|} \simeq 2k_0^2 \quad (7)$$

$$K_{BFKL}^m = \frac{2k_0^2 m^2}{|k_0^2 - m^2|} \simeq 2k_0^2 \quad (8)$$

and convolution with collinear approximated BFKL kernels gives:

$$(K_{2 \rightarrow 2} K_{2 \rightarrow 2}) \otimes \mathcal{V}_{rl\{a'\}}^{LO N_c} = \tilde{\delta}_{a'_1 a'_2} 16\pi k^2 \frac{g^4}{2} \int_{k_0^2}^{k^2} \frac{d^2 l}{(2\pi)^3} \int_{k_0^2}^{k^2} \frac{d^2 m}{(2\pi)^3} \frac{2k_0^2}{l^4} \frac{2k_0^2}{m^4} \frac{(2(l \cdot m)^2 - l^2 m^2)}{k^4} = 0 \quad (9)$$

In collinear approximation we **do not get** contribution from TPV. Leading twist vanishes → **collinear pole vanishes**. Mismach with **GLR**.

## Disconnected part

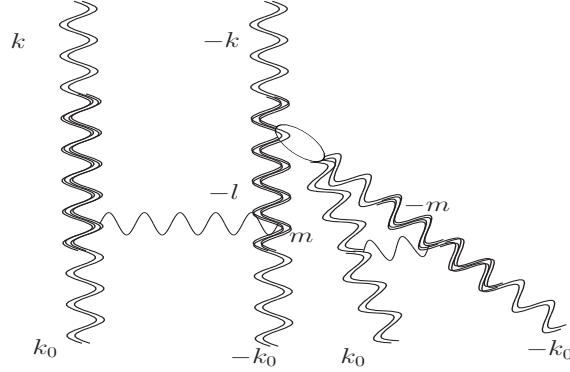


Figure 4:

$G_2$  function reads:

$$G_2(l, m) = -2k^4 \ln \frac{|l|}{|l + m|} \delta(l^2 - k^2) - 2k^4 \ln \frac{|m|}{|l + m|} \delta(m^2 - k^2) \quad (10)$$

Convoluting with impact factor and adding rungs we get:

$$\phi \otimes K_{BFKL} K_{BFKL} \otimes \mathcal{V} \simeq \frac{k_0^6}{Q_1^6} \log \frac{k_0^6}{Q_1^6} \quad (11)$$

↑

negligible in DIS limit

## Anticollinear limit

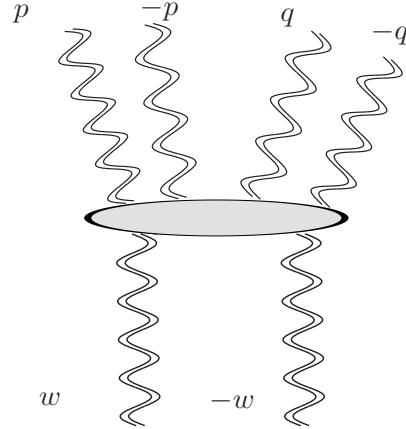


Figure 5:

$$G_1(a, b) = w^2 \left[ \frac{1}{(1 - \frac{w}{a})^2} + \frac{1}{(1 + \frac{w}{b})^2} - \frac{(\frac{w}{a} + \frac{w}{b})^2}{(1 - \frac{w}{a})^2(1 + \frac{w}{b})^2} \right] \quad (12)$$

$w \ll p, q$  the leading term is:

$$\mathcal{V}_{rl\{b'\}}^{LO N_c}(p, -p, q, -q) \simeq \tilde{\delta}_{b'_1 b'_2} 2\pi \frac{g^4}{2} 2w^2 \quad (13)$$

↑

corresponds to the **anticollinear** pole in Mellin space

Comments:

- We do not get logarithmic contribution after convolution with BFKL kernel.
- One gets expected logarithms in when one allows non forward configuration.
- Disconnected pieces do not give expected fourth power of  $\log$

### 3. Angular averaged TPV in BK written for unintegrated gluon density

BK $\leftarrow$  fan diagram equation(Bartels,Lipatov,Vacca)

Averaging over angles yields:

↓

$$\begin{aligned} \mathcal{V}_{\{a\}}^{L0N_c}(a, -a, b, -b) &= \tilde{\delta}_{a'_1 a'_2} \frac{g^4}{2} (2\pi)^2 \left[ 8k^2 \theta(a^2 - k^2) \theta(b^2 - k^2) \right. \\ &\quad \left. + \frac{1}{2} k^4 \ln \left( \frac{a^2}{b^2} \right) \delta(a^2 - k^2) \theta(b^2 - a^2) + \frac{1}{2} k^4 \ln \left( \frac{b^2}{a^2} \right) \delta(b^2 - k^2) \theta(a^2 - b^2) \right] \end{aligned}$$

It's action on Green function  $f^{(2)}(a, b) = f^{(1)}(a)f^{(1)}(b)$ :

↓

$$\begin{aligned} \mathcal{V}_{\{a\}}^{L0N_c}(a, -a, b, -b) \otimes f^{(2)}(a, b) &= C \tilde{\delta}_{a'_1 a'_2} \left( k^2 \int_{k^2}^{\infty} \frac{da^2}{a^4} f(a^2) \int_{k^2}^{\infty} \frac{db^2}{b^4} f(b^2) + \right. \\ &\quad \left. + \frac{1}{2} f(k^2) \int_{k^2}^{\infty} \frac{da^2}{a^4} \ln \left( \frac{a^2}{k^2} \right) f(a^2) + \frac{1}{2} f(k^2) \int_{k^2}^{\infty} \frac{db^2}{b^4} \ln \left( \frac{b^2}{k^2} \right) f(b^2) \right) \end{aligned}$$

Comments:

- main contribution from **anticollinear pole**
- the lower  $k$  the longer path of integration, and stronger nonlinear term
- presence of disconnected contribution

The same structure in BK equation for unintegrated gluon density:  
 (Kimber,Martin,Kwieciński),(Kwieciński,K.K)

$$\frac{\partial f(x, k^2)}{\partial \ln 1/x} = \frac{N_c \alpha_s}{\pi} K \otimes f(x, k^2) \quad (14)$$

$$-\frac{2\alpha_s^2}{R^2} \left[ \textcolor{red}{k}^2 \left( \int_{k^2}^{\infty} \frac{dk'^2}{k'^4} f(x, k'^2) \right)^2 + f(x, k^2) \int_{k^2}^{\infty} \frac{dk'^2}{k'^4} \ln \left( \frac{k'^2}{k^2} \right) f(x, k'^2) \right]$$

related through

$$f(x, k^2) = \frac{N_c}{4\alpha_s \pi^2} k^4 \nabla_k^2 \phi(x, k^2)$$

to:

$$\frac{\partial \phi(x, k^2)}{\partial \ln 1/x} = \bar{\alpha} \pi K \otimes \phi(x, k^2) - \bar{\alpha} \phi^2(x, k^2) \quad (15)$$

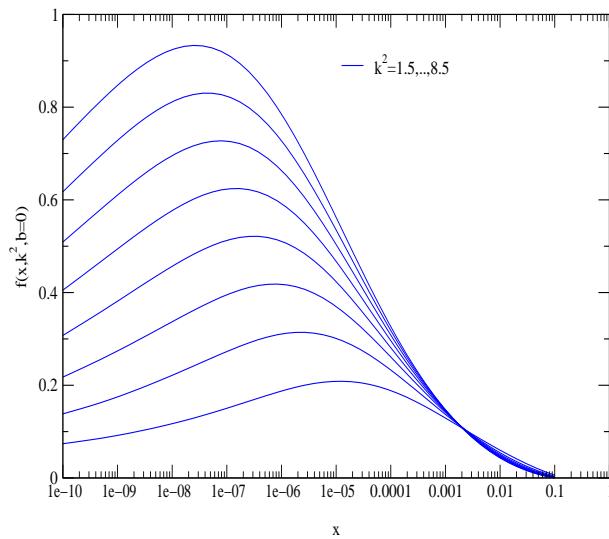


Figure 6:

## Conclusions and additional comments

- $2 \rightarrow 4$  Vertex projected onto pomerons doesn't give contribution in colinear and large  $N_c$  limit
- main contribution comes from anticolinear pole
- collinear poles contribute at finite  $N_c$
- disconnected diagrams give contribution at finite  $N_c$